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Evolution of magnetic states in frustrated diamond lattice antiferromagnetic $\text{Co}(\text{Al}_{1-x}\text{Co}_x)_2\text{O}_4$ spinels

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Using neutron powder diffraction and Monte Carlo simulations we show that a spin-liquid regime emerges at *all* compositions in the diamond-lattice antiferromagnets $\text{Co}(\text{Al}_{1-x}\text{Co}_x)_2\text{O}_4$. This spin-liquid regime induced by frustration due to the second-neighbor exchange coupling J_2 is gradually superseded by antiferromagnetic collinear long-range order ($\mathbf{k}=0$) at low temperatures. Upon substitution of Al^{3+} by Co^{3+} in the octahedral B site the temperature range occupied by the spin-liquid regime narrows and T_N increases. To explain the experimental observations we considered magnetic anisotropy D or third-neighbor exchange coupling J_3 as degeneracy-breaking perturbations. We conclude that $\text{Co}(\text{Al}_{1-x}\text{Co}_x)_2\text{O}_4$ is below the theoretical critical point $J_2/J_1=1/8$, and that magnetic anisotropy assists in selecting a collinear long-range ordered ground state, which becomes more stable with increasing x due to a higher efficiency of O- Co^{3+} -O as an interaction path compared to O- Al^{3+} -O.

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I. INTRODUCTION

Magnetic systems with frustration induced by competing exchange interactions quite often manifest unconventional ground states, the most intriguing of which are spin liquids.¹ Recently one such exotic state, a “spiral spin liquid,” was uncovered theoretically in a classical treatment of diamond-lattice Heisenberg antiferromagnets (AFM) by Bergman *et al.*² who showed that competition between nearest and next-nearest-neighbor exchange couplings J_1 and J_2 creates—for $J_2/J_1 > 1/8$ —a highly degenerate ground state consisting of a set of coplanar spirals, whose propagation vectors form a continuous surface in momentum space. The frustration results in a rich phase diagram as a function of the ratio J_2/J_1 . The degeneracy of these ground states can be lifted by thermal² or quantum³ fluctuations leading to an “order-by-disorder” phase transition from a spiral spin-liquid regime to an ordered state.

Among the diamond-lattice AFM, compounds with the spinel structure recently attracted much attention.^{4–6} In particular, Co-Al oxides were considered as promising candidates for study of order-by-disorder physics.^{2,3} In these compounds of general stoichiometry AB_2O_4 the tetrahedral A sites are occupied by high-spin ($S=3/2$) magnetic Co^{2+} ions which form a diamond lattice consisting of two interpenetrating face-centered cubic sublattices coupled antiferromagnetically. The octahedral B sites can be filled either by nonmagnetic Al^{3+} ions and/or by low-spin ($S=0$) nonmagnetic Co^{3+} ions.

Existing experiments on Co-Al oxide spinels do not provide a clear picture. An early neutron-diffraction study⁷ on Co_3O_4 showed that the magnetic moments of the Co^{2+} ions located at the tetrahedral sites form a simple collinear AFM below the Neel temperature $T_N=40$ K. This picture has been

questioned by a recent muon spin rotation and relaxation (μSR) study⁸ which found two frequency components near T_N suggesting incommensurate magnetic order. Experimental observations on CoAl_2O_4 are also contradictory. A powder neutron-diffraction study of Roth⁹ suggested long-range AFM order below 4 K, while Krimmel *et al.*¹⁰ detected a spin-liquid (or glassylike) ground state. Electron spin resonance, magnetization and specific-heat measurements identified the ground state of CoAl_2O_4 as spin glasslike with a high-frustration parameter $|T_{CW}|/T_N$ of 22 (Ref. 11) and 10 (Ref. 12), respectively. These experimental results are consistent with the calculations of Bergman *et al.*² which placed CoAl_2O_4 in the region of $J_2/J_1 \approx 1/8$ where the spiral surface begins to develop. To clarify the situation, it is essential to understand whether the ground state of Co_3O_4 is a simple collinear antiferromagnet and if CoAl_2O_4 is a spiral spin liquid fluctuating among degenerate spirals. It is also important to understand why substitution within the nonmagnetic B site changes the magnetic properties so drastically. Tristan *et al.*¹³ attempted to answer these questions based on bulk macroscopic measurements of $\text{Co}(\text{Al}_{1-x}\text{Co}_x)_2\text{O}_4$ system. They proposed that the spin-liquid state is realized in CoAl_2O_4 ($x=0$), while with increasing Co substitution x the second-neighbor coupling J_2 decreases and collinear AFM long-range order develops.

In this work we address these important issues and refine the evolution of the magnetic states in $\text{Co}(\text{Al}_{1-x}\text{Co}_x)_2\text{O}_4$ spinels by means of neutron powder diffraction supported by Monte Carlo simulations.¹⁴ Our analysis reveals that for $x=0$ the system is close to the critical point $J_2/J_1=1/8$ in the phase diagram while frustration is weaker for $x>0$. We argue that magnetic anisotropy assists in the selection of a collinear long-range ordered ground state. The difference in T_N and in the extent of the spin-liquid regime apparently

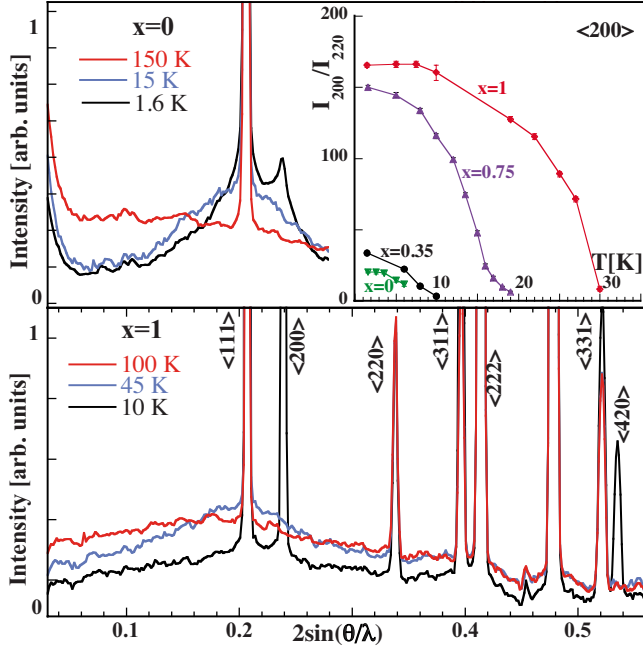


FIG. 1. (Color online) DMC patterns of $x=0$ and 1 samples at selected temperatures. Reasonable statistics on diffuse scattering has been obtained by measuring for about 10 h at one temperature. Inset: temperature dependence of intensity of the $\langle 200 \rangle$ magnetic reflection normalized to the nuclear $\langle 220 \rangle$ intensity.

originates from the decrease in the J_2/J_1 ratio with increasing x .

II. EXPERIMENTAL

Polycrystalline $\text{Co}(\text{Al}_{1-x}\text{Co}_x)_2\text{O}_4$ samples with $x=0, 0.35, 0.75$, and 1 have been prepared as reported in Ref. 11. Thorough structural characterizations by high-resolution x-ray synchrotron diffraction (on the Materials Science beamline at Swiss Light Source, $\lambda=0.41414$ Å) and neutron diffraction (on the HRPT diffractometer at the Swiss Neutron Spallation Source SINQ, $\lambda=1.1545$ Å) indicated the absence of inversion in all samples excepting CoAl_2O_4 . For the latter compound the inversion is 17% and this might influence T_N . Analyzing the lattice constants and x-ray peak profiles we concluded that the Co^{3+} ions are homogeneously and randomly distributed over the B sites. This follows because (i) the lattice constants obey Vegard's law as the function of x and (ii) there is no peak asymmetry, which would have been present in the high-resolution data if any significant degree of inhomogeneity was present.

Medium-resolution neutron powder diffraction patterns for four compositions have been collected in the temperature range 1.5–150 K on the DMC instrument at SINQ with a neutron wavelength $\lambda=2.4526$ Å. For all compositions, broad bumps resembling a liquidlike structure factor and reflecting short-range magnetic order start to develop below the Curie-Weiss temperature $|\theta_{CW}|=110$ K;¹³ they narrow and shift to higher $\sin \theta/\lambda$ with cooling (Fig. 1). Approaching T_N the diffuse scattering localizes near the $\langle 111 \rangle$ and $\langle 200 \rangle$ positions. Below T_N the liquidlike features remain but

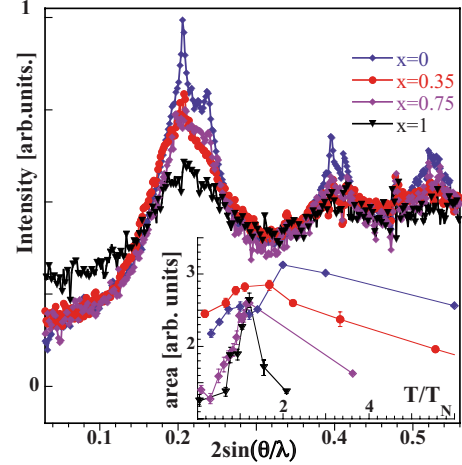


FIG. 2. (Color online) Maximal diffuse magnetic scattering in the differential $T=150$ K neutron powder patterns. The T value is equal to 6, 15, 19, and 45 K for the compositions $x=0, 0.35, 0.75$, and 1, respectively. Inset: evolution of diffuse scattering quantified as an area of Lorentzian fitted to the first bump.

gradually loose spectral weight as magnetic Bragg peaks due to long-range order develop. The spectral weight of the diffuse scattering component continuously increases from Co_3O_4 to CoAl_2O_4 (Fig. 2) as the frustration parameter $|T_{CW}|/T_N$ grows. We can quantify it by the area of a Lorentzian fitted to the first diffuse bump. As shown in the inset of Fig. 2 this area is largest roughly at $T/T_N \approx 1$ for all compositions, but for $x=0$ and 0.35 diffuse scattering develops far above and remains significantly below this value. For $x=1$ the T/T_N interval revealing diffuse scattering is narrower but still much too large to be interpreted as classical critical scattering of a three-dimensional long-range ordered magnet.

Regarding the long-range order, all samples show magnetic Bragg peaks at low temperatures. The ordering temperature T_N and the static ordered magnetic moment decrease from Co_3O_4 to CoAl_2O_4 (see Table I and inset of Fig. 1). It should be noted that the $x=0$ diffraction pattern does not correspond to a conventional long-range ordered state: diffuse scattering clearly dominates and the $\langle 200 \rangle$ peak is so broad and weak that the ordered moment cannot be determined with high accuracy. Nevertheless, the magnetic Bragg pattern is the same for all compositions and is consistent with the collinear two-sublattice model proposed by Roth.⁷

III. MODELING

To model the observed magnetic diffuse scattering we used the quasistatic approximation,¹⁵ which assumes that the

TABLE I. The Néel temperature, T_N , and the ordered magnetic moment, M , determined at 1.6 K from the DMC patterns.

x	T_N (K)	M (μ_B)
1.0	29	3.53(3)
0.75	16.5	2.59(3)
0.35	9	1.31(4)
0.0	5	0.25(7)

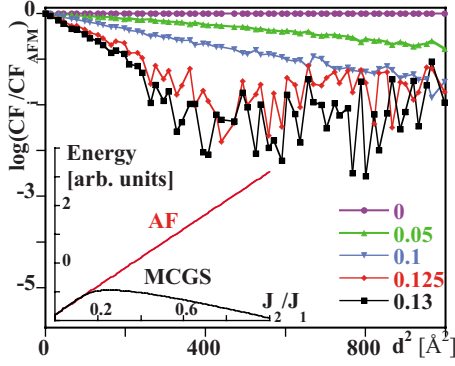


FIG. 3. (Color online) Log of ratio of spin-pair correlation functions $|CF_i/CF_{AFM}|$ versus d^2 for selected J_2/J_1 ratios. Inset: the minimal energy curve for a collinear antiferromagnet (AF, red) and for the best Monte Carlo ground state (MCGS, black).

observed diffuse scattering can be attributed to static correlations. This approach is justified since the bandwidth of magnetic excitations is below 6 meV whereas all scattering with energies below 14.7 meV is summed in our diffraction experiment. We performed a Monte Carlo search for the ground state of a cluster of 1047 atoms. The energy of the classical Heisenberg Hamiltonian for spins \mathbf{S} interacting with first- and second-neighbor antiferromagnetic couplings $J_1, J_2 > 0$

$$H = J_1 \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j + J_2 \sum_{\langle\langle ij \rangle\rangle} \mathbf{S}_i \cdot \mathbf{S}_j \quad (1)$$

was minimized ($\langle\langle ij \rangle\rangle$ means that the sum runs over first-neighbor pairs, $\langle\langle ij \rangle\rangle$ on second-neighbors pairs). If we add either a third-neighbor interaction or an anisotropy contribution, the Hamiltonian will be completed with an additional term ΔH , given, respectively, as

$$\Delta H = J_3 \sum_{\langle\langle\langle ij \rangle\rangle\rangle} \mathbf{S}_i \cdot \mathbf{S}_j \quad (\text{third neighbor}); \quad (2)$$

$$\Delta H = D \sum_i (\mathbf{S}_i \cdot \mathbf{u})^2 \quad (\text{anisotropy}), \quad (3)$$

where $\langle\langle\langle ij \rangle\rangle\rangle$ refers to third-neighbors pairs and \mathbf{u} is the anisotropy direction, that we took as (111). D is the magnitude of the anisotropy term. The moments were kept equal and constant in magnitude; their direction was changed at random, one at the time, and to obtain the ground state only energy-decreasing moves were accepted in the final stage. The stopping criterion for the Monte Carlo was that the last 1000 accepted configurations had an energy spread less than $10^{-6}J_1$. Runs have been repeated for the whole range $0 \leq J_2/J_1 \leq 1$. We calculated the static spin-pair correlation functions and diffraction patterns for the Monte Carlo ground state (MCGS) and the reference AFM clusters for each J_2/J_1 . The correlation function is given by

$$CF(d) = \frac{2}{3} \sum_{ij} \mathbf{S}_i \cdot \mathbf{S}_j \delta(|r_i - r_j| - d), \quad (4)$$

where d is the distance between spins at positions r_i and r_j . We found that for the range $0 < J_2/J_1 < 1/8$ the collinear

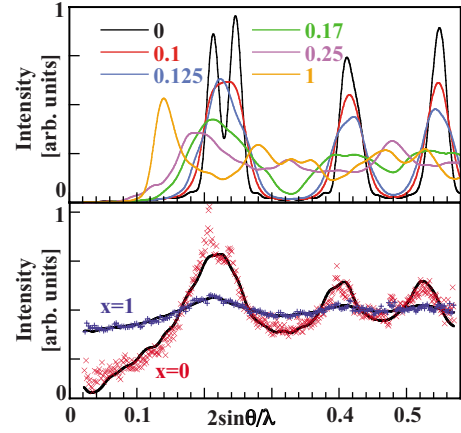


FIG. 4. (Color online) Top: calculated powder diffraction patterns for several selected J_2/J_1 ratios at $T=0$ K. Bottom: observed (symbols) and calculated (line) powder diffraction patterns for $x=0$ and 1 corresponding to the best model parameters. The calculated patterns are corrected for the magnetic form factor of Co^{3+} and polynomial background.

AFM is the ground state, in agreement with Ref. 2. Remarkably, the MCGS was always configurationally different from a collinear AFM state even if the energy difference was negligible (order of 1 mK or less). Analytical calculations (Fig. 3 inset), similar to Ref. 2, support this result. In fact, the energy minimum corresponding to the ground state for $J_2/J_1 < 1/8$ is extremely flat around the $q=0$ point of the first Brillouin zone, and therefore, very many states are in the range of thermal excitation at any accessible temperature.

The MCGS starts closely related to the collinear AFM state, but progressively departs from it with increasing J_2/J_1 . This is clearly seen in the decay of spin correlations, which can be quantified by the ratio CF_{MCGS}/CF_{AFM} . In Fig. 3 we plot $\log|CF_{MCGS}/CF_{AFM}|$ versus squared distance d^2 between spins. Near $J_2/J_1 = 1/8$ the decay changes from predominantly Gaussian $e^{-1/2(d/w)^2}$ (a straight line in this choice of axes) to a mixed Gaussian-exponential form $e^{-d/L} \cdot e^{-1/2(d/w)^2}$ (a curve bending downwards for $200 < d^2 < 400 \text{ \AA}^2$). This behavior can be attributed to the incipient spiral surface. In the calculated diffraction patterns (Fig. 4, top) these changes are evidenced as progressive broadening and distortion of Bragg peaks. Pure Gaussian correlation decay yields a Gaussian peak broadening, whereas Gaussian-exponential decay yields Voigtian point spread (that is, with a sizable Lorentzian convolution component). For $J_2/J_1 > 0.5$ the correlations are completely changed and features of the collinear AFM state fade out.

We now compare the experimental and calculated diffuse scattering. The patterns shown in Fig. 2 were obtained as differences between data measured at $T=6, 15, 19$, and 45 K (for $x=0, 0.35, 0.75$, and 1, respectively) and 150 K. In the calculated patterns in order to obtain the ordered AFM state at $T < T_N$ we needed to perturb the simple Heisenberg Hamiltonian so that the AFM minimum would be deeper and the AFM configuration more stable. For this we considered separately two possible degeneracy-breaking perturbations: magnetic anisotropy and finite J_3 coupling. To retain a simple

model we described anisotropy as a single-ion parameter D . The MC procedure was now changed, allowing for system equilibrations at different temperatures. All moves that would decrease the energy or that would increase it with probability $\propto e^{-\Delta E/T}$ were accepted.¹⁶ When the average and fluctuations of the energy of a large number of the last accepted states were sufficiently stable the temperature was changed. Fixing $J_1=1$ as a convenient energy scale, J_2 , J_3 , D , or T were varied, each in several steps.

Equally good fits were obtained with J_3 and magnetic anisotropy D as perturbations, signifying that the available diffraction data are not sufficient to determine which of these terms stabilizes the ordered ground state. However, we give preference to the magnetic anisotropy following the theoretical study of Ref. 17. The best fits presented in Fig. 4 (bottom) suggest the ratio $J_2/J_1=0.125$ for CoAl_2O_4 and $J_2/J_1=0.05$ for Co_3O_4 . The change in the exchange energy with substitution apparently originates from the peculiarities of the electronic structure. Band-structure analysis¹⁸ shows that near the Fermi level in Co_3O_4 there are Co^{3+} d and oxygen p states, while in CoAl_2O_4 the Al p states are absent from the Fermi level and the weight of O p is diminished. This implies that the interaction path O- Co^{3+} -O is more effective and the corresponding exchange integrals are larger in Co_3O_4 .

Our findings give a natural explanation of the μSR results⁸ on Co_3O_4 as due to spin-liquid physics rather than incommensurate magnetic order. Also, based on the position of CoAl_2O_4 in the J_2/J_1 phase diagram, we suggest that even in an ideal sample, with no inversion or other perturbation, the ground state would be a collinear AFM, though one might need very low temperatures to reach it. As the ground state is the collinear AFM the general term “spin liquid” and not “chiral spin liquid” is appropriate for the high-

temperature regime above T_N in the title system.

A single-crystal inelastic neutron-scattering experiment would allow to extract the absolute values of J_1, J_2 and to validate our conclusions. We remark that the elaborated approach to fit measured diffuse magnetic neutron scattering to Monte Carlo simulations can be easily adapted to other frustrated systems and would be useful in justification of an anticipated Hamiltonian.

IV. SUMMARY

In summary we studied the evolution of magnetic states in $\text{Co}(\text{Al}_{1-x}\text{Co}_x)_2\text{O}_4$ polycrystalline samples with temperature and substitution in the B site. We observed short-range and long-range orders for all compositions. Employing Monte Carlo simulations we found that for $x=0$ the system is in the vicinity of the critical point $J_2/J_1=1/8$, where the spiral spin liquid develops,² but stays in the weakly frustrated limit for $x>0$. We also found that replacement in the nonmagnetic B site changes the strength of exchange interactions which, in turn, leads to significant differences in the ordering temperatures and in the extent of the spin-liquid regime.

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¹Spin liquid is characterized by rapid (at least exponential) decay of spin-spin correlations and the correlation length does not exceed the interatomic distance (Ref. 19), which result in a spin liquidlike structure factor.

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